Asymmetric aggregation enables red-light CO_2 reduction with tunable activity and selectivity by intermolecular electronic coupling

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ABSTRACT

Photocatalytic CO_2 reduction to value-added chemicals is a promising way to simultaneously mitigate environmental and energy issues. Integrating low-energy red light harvesting, fast charge separation, and robust CO_2 activation in a photosystem is a crucial prerequisite for highly efficient CO_2 photoreduction but remains a huge challenge. Herein, we constructed an aggregated [Ru(2, 2'-bipyridine)$_3$]Cl$_2$ (denoted as Ru(bpy)$_3$Cl$_2$) photosystem. Various characterizations combining molecular dynamic simulations and theoretical calculations confirm that the electronic coupling between adjacent Ru(bpy)$_3$Cl$_2$ molecules can induce orbital hybridization, thus broadening the light absorption edge to the red light region. Meanwhile, the asymmetric aggregation of Ru(bpy)$_3$Cl$_2$ enables exciton dissociation and charge transfer through a dipole polarization effect. Significantly, this dipole polarization can also upshift the d-orbital of Ni catalytic centers, greatly strengthening the activation of CO_2 and lowering the formation energy barrier of COOH* intermediates. Consequently, the apparent quantum yield (AQY) of aggregates for CO_2-to-CO reduction reaches 4.2% at 610 nm, significantly higher than the AQYs of other reported photosystems at wavelengths $\geq$ 600 nm. More importantly, the catalytic efficiency and selectivity of CO_2 reduction can be rationally tuned by controlling the aggregated degrees of Ru(bpy)$_3$Cl$_2$. This work highlights the key role of Ru(bpy)$_3$Cl$_2$ aggregation in prompting red-light harvesting and CO_2 activation, which may help boost CO_2 photoreduction efficiency from a fresh perspective.

1. Introduction

Photocatalytic CO_2 reduction into value-added products is a promising path to address energy and environmental issues [1–3]. In the solar spectrum, visible and near-infrared (NIR) light account for 95% of the solar spectrum [4]. Given solar energy utilization, photocatalysts need to capture the broad spectrum of solar light. Still, most existing photocatalysts are usually active under ultraviolet or narrow visible light and display serious charge recombination [5]. Especially, many organic photocatalysts have poor stability because high-energy ultraviolet and visible light irradiation could cause structural transformation [6]. In principle, low-energy red/NIR-light-responsive photosystems may overcome these issues. However, only a few photocatalysts can drive CO_2 reduction with red or NIR light. For example, Kang and co-workers constructed a CNDs/Cu$_2$O photocatalytic system, which showed high photocatalytic activities under NIR light irradiation by up-conversion

[7]. Xiong and co-workers employed the Au rod as a plasmonic light-harvesting unit and copper-palladium alloy shell as a co-catalyst, achieving CO_2 reduction under full-spectrum light illumination [8]. Xie and co-workers synthesized an ultrathin oxygen-deficient cubic WO$_3$ layers, realizing IR-driven CO_2 overall splitting by a defect-induced intermediate band [9]. However, these reported systems show moderate photocatalytic activity. This is because most red/NIR-light-response photocatalysts typically exhibit narrow energy gaps, leading to improper band-edge positions for CO_2 reduction [9,10]. Furthermore, most existing photocatalysts have inefficient CO_2 activation owing to the chemical inertness of linear CO_2. Therefore, it is significant to construct a photosystem that possesses red/NIR light harvesting, efficient charge separation, suitable redox potentials, and robust CO_2 activation simultaneously.

The π-conjugated organic molecules are of growing interest in the field of photocatalysis. It has been proven that molecular packing/

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aggregation can control optical properties in organic nanomaterials [11, 12]. For instance, our group observed that the aggregates of light-active structures can significantly broaden their light-harvesting regions, thereby improving photocatalytic performance [13,14]. Furthermore, the unique electronic coupling within aggregates facilitates charge localization [15,16]. Especially, the spontaneous symmetry breaking of aggregates can generate large dipole moments, which undoubtedly accelerates charge separation and transfer [17,18]. It should be noted that this dipole polarization effect may also regulate the electronic properties of co-catalysts, in principle, probably expediting the activation of CO_2 molecules. Based on the above analysis, the aggregation of light-harvesting units can broaden the light-harvesting region and promote the separation and transfer of charge carriers [19,20]. Significantly, the aggregation of light harvesting units may consolidate the activation of CO_2 by regulating the electronic properties of co-catalysts. However, to the best of our knowledge, the aggregation-enhanced red-light CO_2 reduction has not been examined, and the link between the aggregation of light harvesting units and CO_2 reduction activity remains unclear.

In this work, as a proof-of-concept demonstration, we chose [Ru(2, 2'-bipyridine)$_3$]Cl$_2$ (denoted as Ru(bpy)$_3$Cl$_2$) as a light-harvesting unit and systematically studied the aggregated effect of Ru(bpy)$_3$Cl$_2$ on the red-light-driven CO_2 reduction performance. The structures, photophysical properties, and CO_2 photoreduction activity of Ru(bpy)$_3$Cl$_2$ aggregates were analyzed carefully. Results confirm that the electronic coupling between adjacent Ru(bpy)$_3$Cl$_2$ molecules can induce orbital rearrangement, thus extending the light absorption edge to the red light region. Simultaneously, the spontaneous symmetry breaking of Ru(bpy)$_3$Cl$_2$ aggregates generates a large dipole moment, which boosts the separation and transfer of charge carriers and prolongs the lifetime of charge carriers. Interestingly, this dipole polarization of Ru(bpy)$_3$Cl$_2$ aggregates can also upshift the d-orbital of Ni catalytic centers, promoting the activation of CO_2 and lowering the formation energy barrier of COOH* species. As a result, the Ru(bpy)$_3$Cl$_2$ aggregates exhibit excellent CO_2-to-CO reduction activity under red light irradiation with an apparent quantum yield of 4.2% at 610 nm. More importantly, the efficiency and selectivity of this CO_2 photoreduction system can be tuned by changing the aggregated degrees of Ru(bpy)$_3$Cl$_2$ molecules.

2. Experimental

2.1. Chemicals

All chemicals were analytical-grade and used without further purification. Acetonitrile (≥99.9%), triethanolamine (≥99.5%), sodium sulfate anhydrous (99.5%), and 2,2-bipyridine (≥99.5%) were purchased from Sinopharm Chemical Reagent Co., Ltd. 2, 2'-Bipyridine nickel dichloride (≥99.9%) and cetyl trimethyl ammonium bromide (≥99.9%) were purchased from Macklin. [Ru(bpy)$_3$]Cl$_2$·6 H$_2$O (≥99.0%) was obtained from Alfa Aesar.

2.2. Preparation of droplet-confined Ru(bpy)$_3$Cl$_2$

Ru(bpy)$_3$Cl$_2$ with different amounts were respectively dissolved in water/acetonitrile (2 mg/mL for droplet-A, 4 mg/mL for droplet-B, 8 mg/mL for droplet-C). Then, the above Ru(bpy)$_3$Cl$_2$ solution was injected into 10 mL hexane containing 30 mg cetyl trimethyl ammonium bromide, while ultrasonic treatment was carried out for 5 min under continuous stirring to form uniform droplets (the overall mass of Ru(bpy)$_3$Cl$_2$ within the above three droplets are controlled to be 8 mg).

2.3. Characterization

UV–vis absorption spectra of all samples were analyzed using a HITACHI UH5700 UV-Vis-NIR spectrophotometer. Photoluminescence (PL) spectra and time-resolved photoluminescence of all samples were measured by a FL spectrophotometer (FLSP920, FLS1000). Confocal laser microscopy images were obtained with a Carl Zeiss LSM880 instrument. Small-angle X-ray scattering (SAXS) experiment was measured on model Xeuss 2.0 instrument from Xenocs, France, the incident beam was Cu Kα radiation (λ = 0.1542 nm). The data was fitted to the model in the SasView software package (version 5.0.6). The error is obtained from the fitting software and does not take into account other sources of error. Cryogenic transmission electron microscopy (Cryo-TEM) was measured by a FEI Talos F200C. The ultrafast femtosecond transient absorption spectra (*fs*-TA) measurement was conducted on a Helios spectrometer using a regeneratively amplified femtosecond Ti: sapphire laser system (Spitfire Pro-F1KXP, Spectra-Physics). Diffuse reflectance infrared Fourier transform spectrum was measured by BRUKER, EQUINOX 55. $^{13}CO_2$ isotope labeling experiment was performed under the same reaction conditions. A gas chromatography-mass spectrometry was used to analyze the generated ^{13}CO. Linear scanning voltammetry experiments were measured by CHI660E C22142. The working electrode was carbon paper, platinum wire as the counter electrode, Ag/AgCl as the reference electrode and the electrolyte solution was an aqueous solution of 0.2 M Na$_2$SO$_4$.

2.4. Photocatalytic reaction and apparent quantum yield measurements

Photocatalytic CO_2 reduction reaction was conducted in a closed reaction system. In brief, Ru(bpy)$_3$Cl$_2$ with different amounts were added to 2 mL of water and 3 mL of acetonitrile. 2, 2'-Bipyridine nickel dichloride (Ni(bpy)Cl$_2$) and triethanolamine (TEOA) were used as co-catalyst and electron donor, respectively. The air in the reaction bottle was expelled before the reaction, and then the reaction system was filled with CO_2 gas. A 50 W monochromatic LED lamp (610 nm) was used as the light source. The light intensity is controlled to be 110.2 mW·cm^{-2}. The reaction temperature was maintained at room temperature (30 ± 2°C) by using a cooling fan. The rotational speed of the magnet is maintained at 600 rmp. The products during the photocatalytic reaction were quantified by gas chromatography (TCD detector).

The apparent quantum yield (AQY) of CO_2 reduction was measured by using a monochromatic LED lamp as excitation light. The irradiation area was controlled at 0.196 cm^2. The reaction temperature was controlled at room temperature (30 ± 2 °C). The amount of catalyst is 44.8 mg. The rotational speed of the magnet is maintained at 600 rmp. The light intensity of 610 nm was 26.48 mW·cm^{-2} (ILT 950 spectroradiometer).

$$AQY = Ne/Np*100\% = 2M\ NAhc/SPt\lambda*100\%$$

Where M is the amount of CO (mol), NA is Avogadro constant (6.02×10^{23} mol^{-1}), h is Planck constant (6.626×10^{-34} J • s), c is light speed (3×10^8 m/s), S is irradiation area (cm^2), P is light intensity (W/cm^2), t is reaction time (s), λ is the wavelength of light irradiation (m).

2.5. Computational details

Classical molecular dynamic (MD) simulations were performed to investigate the structures of Ru(bpy)$_3$Cl$_2$ in the GROMACS (version 2021.7) simulation package. All of the initial structures were optimized at B3LYP level in Gaussian 09 software. The simulated parameters of Ru(bpy)$_3$Cl$_2$ were obtained using UFF force filed [21], which has confirmed the accuracy in pieces of literature [22]. During MD simulation, the Particle-mesh Ewald was selected for long-range electrostatics. The cutoff was set at 1.0 nm, and the time step was set at 1 fs. In detail, for pure acetonitrile (ACN) and acetonitrile/cyclohexane (CYH) environments containing the same number of Ru(bpy)$_3$Cl$_2$ molecules at the size of a = 5.0 nm, b = 5.0 nm, c = 5.0 nm, α = β = γ =90°. Therefore, the 300 ns MD simulations were performed for each system. After equilibration checking, only the final 100 ns converged trajectories were used for further analysis.

All calculations were performed by Gaussian 09 software. The B3LYP functional and mixed basis set of SDD were used to calculate the metal atoms, and the nonmetal atoms were used for geometry optimization of the molecule using 6–31 G* [23,24]. The convergence accuracy of the SCF reaches the default convergence standard. Subsequently, harmonic frequency analysis calculations at the same level were carried out to indicate the absence of imaginary frequencies at the standing point. To obtain more accurate energy, single-point calculation was performed using the higher basis set, where 6–311 G* was used for nonmetal atoms and SDD for metal atoms in a mixture of water and acetonitrile solvents, using the SMD solvent model [25]. The Gibbs free energy was then corrected to the refining energy at 298.15 K and 1 atm. In this study, we discuss the reaction mechanism in terms of the corrected Gibbs free energy given under the corresponding compounds.

Moreover, we calculated electron-hole distribution, excited states, electron density difference, and Mulliken charge populations of the molecules. The geometrical configuration, highest occupied molecular orbital (HOMO)-lowest unoccupied molecular orbital (LUMO) energy gap, excited state orbital distribution, the electrostatic potential at the molecular surface, and dipole moments of $Ru(bpy)_3Cl_2$ have also been visualized by using Multiwfn 3.8 [26] and VMD 1.9.3 [27] software.

3. Results and discussion

3.1. Structural and characterization of $Ru(bpy)_3Cl_2$ aggregates

Fig. 1a-b display the UV–vis absorption spectra of $Ru(bpy)_3Cl_2$ with different concentrations. The $Ru(bpy)_3Cl_2$ monomer shows a characteristic peak at 450 nm [28,29]. With the increase of $Ru(bpy)_3Cl_2$ concentrations, the characteristic peak at 450 nm is noticeably red-shifted and the light absorption intensity from 500 to 700 nm can be enhanced significantly [30,31]. This result indicates the formation of $Ru(bpy)_3Cl_2$ aggregates, which is further supported by the calculated bandgaps result of $Ru(bpy)_3Cl_2$ aggregates (Figure S1) [32,33]. To further investigate the formation of $Ru(bpy)_3Cl_2$ aggregates, we analyzed the photoluminescence (PL) spectra of $Ru(bpy)_3Cl_2$ with

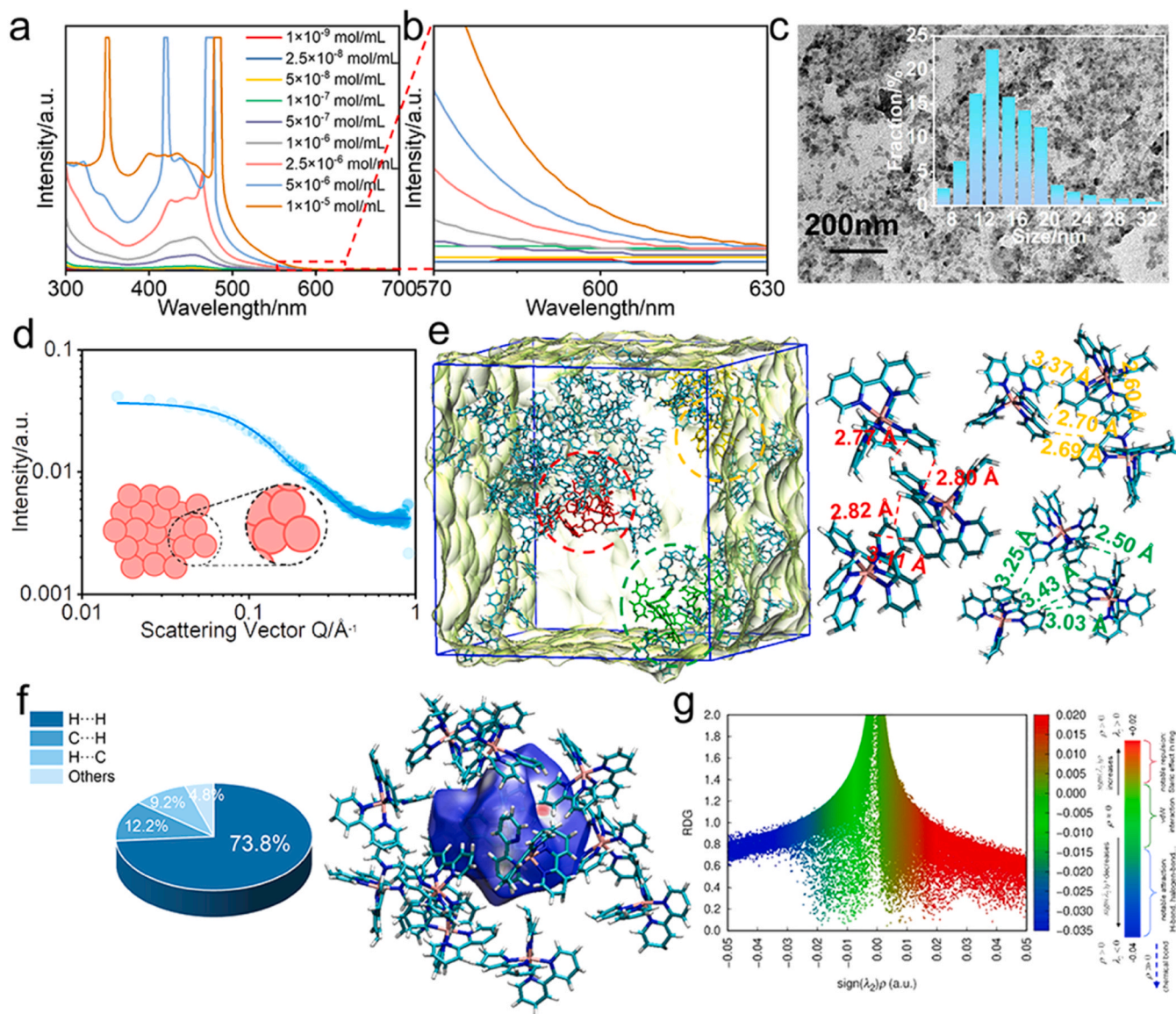**Fig. 1.** (a-b) UV–vis absorption spectra of $Ru(bpy)_3Cl_2$ with different concentrations, (c) Cryo-TEM image of $Ru(bpy)_3Cl_2$ aggregates, (d) SAXS data of $Ru(bpy)_3Cl_2$ aggregates, (e) The structures of $Ru(bpy)_3Cl_2$ aggregates by MD simulation, (f) Hirschfeld surface analysis (mapped over d_{norm}), and proportions of intermolecular C⋯H, H⋯H, H⋯C, and other interactions to the total intermolecular interactions based on MD simulation, and (g) The function of reduced density gradient and sign $(\lambda_2)\rho$ scatter spectra for $Ru(bpy)_3Cl_2$ aggregates based on our MD simulation.

different concentrations (Figure S2). It is obvious that the PL emission band is red-shifted with increasing the concentrations of Ru(bpy)$_3$Cl$_2$. This red-shift behavior of PL emission also reflects the formation of aggregates. The formation of Ru(bpy)$_3$Cl$_2$ aggregates was imaged using cryogenic transmission electron microscopy (Cryo-TEM) to capture the solution-state morphology and size. As shown in Fig. 1c, the Ru(bpy)$_3$Cl$_2$ aggregates display sphere-like morphology with a narrow size distribution. The small-angle X-ray scattering (SAXS) curve was further employed to determine the structures of Ru(bpy)$_3$Cl$_2$ aggregates (Fig. 1d). The fitting data indicates the existence of sphere-like nanostructures [34,35]. The vanishing of the linear Guinier region can be attributed to aggregation in the sample (Figure S3) [35–37], agreeing with the Cryo-TEM measurement.

To further analyze the microstructures of Ru(bpy)$_3$Cl$_2$ aggregates, a molecular dynamic (MD) simulation was performed (Fig. 1e). The Ru(bpy)$_3$Cl$_2$-Ru(bpy)$_3$Cl$_2$ contact map authorizes that the Ru(bpy)$_3$Cl$_2$ molecules tend to interact with adjacent Ru(bpy)$_3$Cl$_2$ molecules (Figure S4). The distances of neighboring Ru(bpy)$_3$Cl$_2$ molecules are ranged from 2.5 to 3.59 Å (Figure S5). Interestingly, the distances and angles of adjacent Ru(bpy)$_3$Cl$_2$ molecules are different, suggesting the asymmetric feature of Ru(bpy)$_3$Cl$_2$ aggregates. Therefore, Hirshfeld surface analysis based on MD simulation was performed to quantitatively estimate the intermolecular interactions (Fig. 1f and S6) [38,39]. The weak intermolecular C⋯H, H⋯H, and H⋯C interactions account for more than 95% of total intermolecular interactions, indicating that weak non-covalent interactions mainly contribute to the aggregation of Ru(bpy)$_3$Cl$_2$. Reduced density gradient (RDG) scattering spectrum also ratifies the presence of weak intermolecular interactions in the Ru(bpy)$_3$Cl$_2$ aggregates (Fig. 1g), consistent with the findings of Hirshfeld surface analysis [40,41].

Considering the red-shifted light absorption of Ru(bpy)$_3$Cl$_2$ aggregates, the Ru(bpy)$_3$Cl$_2$-Ru(bpy)$_3$Cl$_2$ orbital interactions analysis is necessary to understand the electronic coupling between adjacent Ru(bpy)$_3$Cl$_2$ molecules [42,43]. As revealed in Fig. 2, the LUMO orbital of the trimer mainly originates from the 131st orbital of Fragment 1 (95%), with a decrease in energy of 5.87 eV. The HOMO orbital mainly originates from the 130th orbital of Fragment 3 (94%), with an increase in energy of 5.11 eV. The LUMO+1 and HOMO-1 orbitals are respectively derived from the 134th orbital of Fragment 1 (64%) and the 130th orbital of Fragment 2 (75%). These results confirm the existence of strong electronic coupling between adjacent Ru(bpy)$_3$Cl$_2$ molecules, which rearranges their molecular orbitals (MOs), in agreement with the orbital analysis of the dimer and tetramer (Figures S7 and S8) [44]. The above results suggest intermolecular weak interactions can induce orbital rearrangements and hybridization of Ru(bpy)$_3$Cl$_2$ molecules, thus generating new light-harvesting properties. It should be noted that weak non-covalent interactions are not unusual in organic molecules but have been overlooked in the past. The above discovery will deepen the

understanding of intermolecular weak forces and offer a promising route to regulate the physicochemical properties, especially the optical properties of photocatalytic materials from a fresh perspective.

In principle, the asymmetric feature of Ru(bpy)$_3$Cl$_2$ aggregates could induce the generation of large dipole moments, which boosts the separation of photogenerated charge carriers. Therefore, we statistically analyzed the dipole moments of Ru(bpy)$_3$Cl$_2$ aggregates with different aggregated degrees based on MD simulation. The dipole moment of Ru(bpy)$_3$Cl$_2$ monomer is calculated to be 0.509 Debye (denoted as 0.509 D) (Figure S9). Significantly, the dipole moments of aggregates increase with increasing the aggregated degrees of Ru(bpy)$_3$Cl$_2$ (Figures S9-12). The large dipole moments of the Ru(bpy)$_3$Cl$_2$ aggregates can be primarily ascribed to structural symmetry breaking [45]. Interestingly, we also observed the dynamical behaviors of Ru(bpy)$_3$Cl$_2$ aggregates due to the weak interactions between adjacent Ru(bpy)$_3$Cl$_2$ molecules. Randomly selected snapshot images of Ru(bpy)$_3$Cl$_2$ aggregates recorded at 94, 95, 96, 97, 98, and 99 ns display different local structures (Fig. 3) and electron density distributions (Figure S13), reflecting their dynamical nature. This dynamic behavior was found to stabilize the asymmetric structures of Ru(bpy)$_3$Cl$_2$ aggregates. It has been widely accepted that large dipole moments can result in a strong built-in electric field, which will certainly accelerate exciton dissociation and charge transfer [46,47]. To confirm this deduction, we calculated the exciton binding energies of Ru(bpy)$_3$Cl$_2$ aggregates by using the Multiwfn program. Impressively, the exciton binding energy sharply decreases from 4.767 eV for Ru(bpy)$_3$Cl$_2$ monomer to 0.730 eV for tetramer aggregates, consistent with the results of electron-hole distributions (Figure S14). The decreased exciton binding energy reflects that the aggregation of Ru(bpy)$_3$Cl$_2$ can expedite exciton dissociation into free electrons and holes, which increases the probability of electrons and holes participating in redox reactions [48,49].

Furthermore, the steady-state PL emission intensity of Ru(bpy)$_3$Cl$_2$ aggregates is lower than that of Ru(bpy)$_3$Cl$_2$ monomer, which suggests an enhanced charge separation of Ru(bpy)$_3$Cl$_2$ aggregates (Figure S15). Also, we observed that the PL lifetime of Ru(bpy)$_3$Cl$_2$ aggregates can be significantly prolonged (Figure S15), signifying an enhanced chance for photogenerated electrons participating in CO$_2$ reduction, in agreement with the steady-state PL emission results. To deeply explore the aggregated effect of Ru(bpy)$_3$Cl$_2$ on charge separation and transfer dynamics, we estimated the femtosecond transient absorption (*fs*-TA) spectra of Ru(bpy)$_3$Cl$_2$ monomer and aggregates, as illustrated in Figs. S16 and 4b. *fs*-TA demonstrates the absorption spectra versus time and wavelength. The Ru(bpy)$_3$Cl$_2$ monomer displays a weak absorption (excited state absorption, ESA) [50,51]. The fitting lifetime at 615 nm is 3.86 ns (Table S2). Whereas, the Ru(bpy)$_3$Cl$_2$ aggregates exhibit a strong ESA signal. Compared to the Ru(bpy)$_3$Cl$_2$ monomer, the Ru(bpy)$_3$Cl$_2$ aggregates exhibit different dynamic behavior, which reflects the unique interactions within aggregates. The fitting lifetime at 615 nm is 28.5 ns for Ru(bpy)$_3$Cl$_2$ aggregates, which is longer than that for Ru(bpy)$_3$Cl$_2$ monomer. Significantly, a new transient absorption band at 755 nm was witnessed, which belongs to the charge transfer (CT) state [50,51]. The fitting lifetime of CT state is 18.7 ns (Fig. 4c and Table S2). It is widely recognized that the CT state can easily overcome the Coulombic attraction of excitons to form free charges on the time scale of picoseconds. The formation of CT state reflects an effective exciton dissociation and charge transfer dynamics of Ru(bpy)$_3$Cl$_2$ aggregates, consistent with the above results of DFT predictions, steady-state PL, and time-resolved PL spectra [52–55]. Moreover, we have also estimated the CO$_2$ absorption ability of the aggregates (Figure S17). The aggregates show obvious CO$_2$ absorption capability, providing the possibility for CO$_2$ reduction.

Fig. 2. Orbital interaction diagram for Ru(bpy)$_3$Cl$_2$ aggregates. The vertical axis shows MO energies in eV. Blue solid and dashed bars correspond to occupied and unoccupied MOs, respectively. The MOs are plotted as isosurface graphs with isovalue of 0.02 a.u.

3.2. Photocatalytic activity test

Intrigued by the enhanced light harvesting and the effective charge separation of Ru(bpy)$_3$Cl$_2$ aggregates, we investigated their

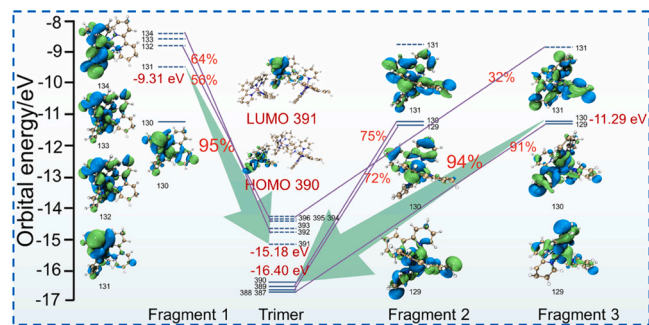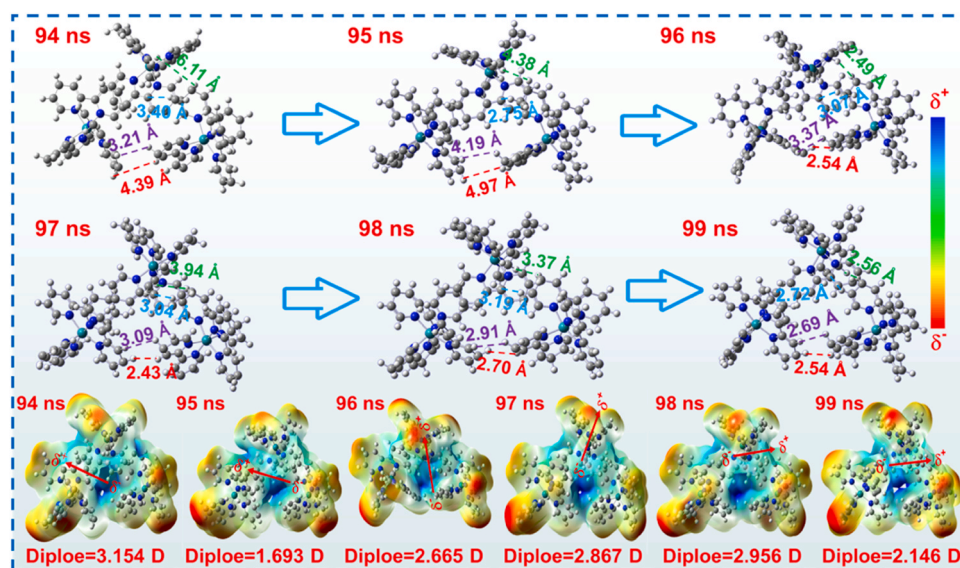**Fig. 3.** Dynamic properties of trimer aggregates in MD snapshot with related dipole moments and electrostatic potential distributions.
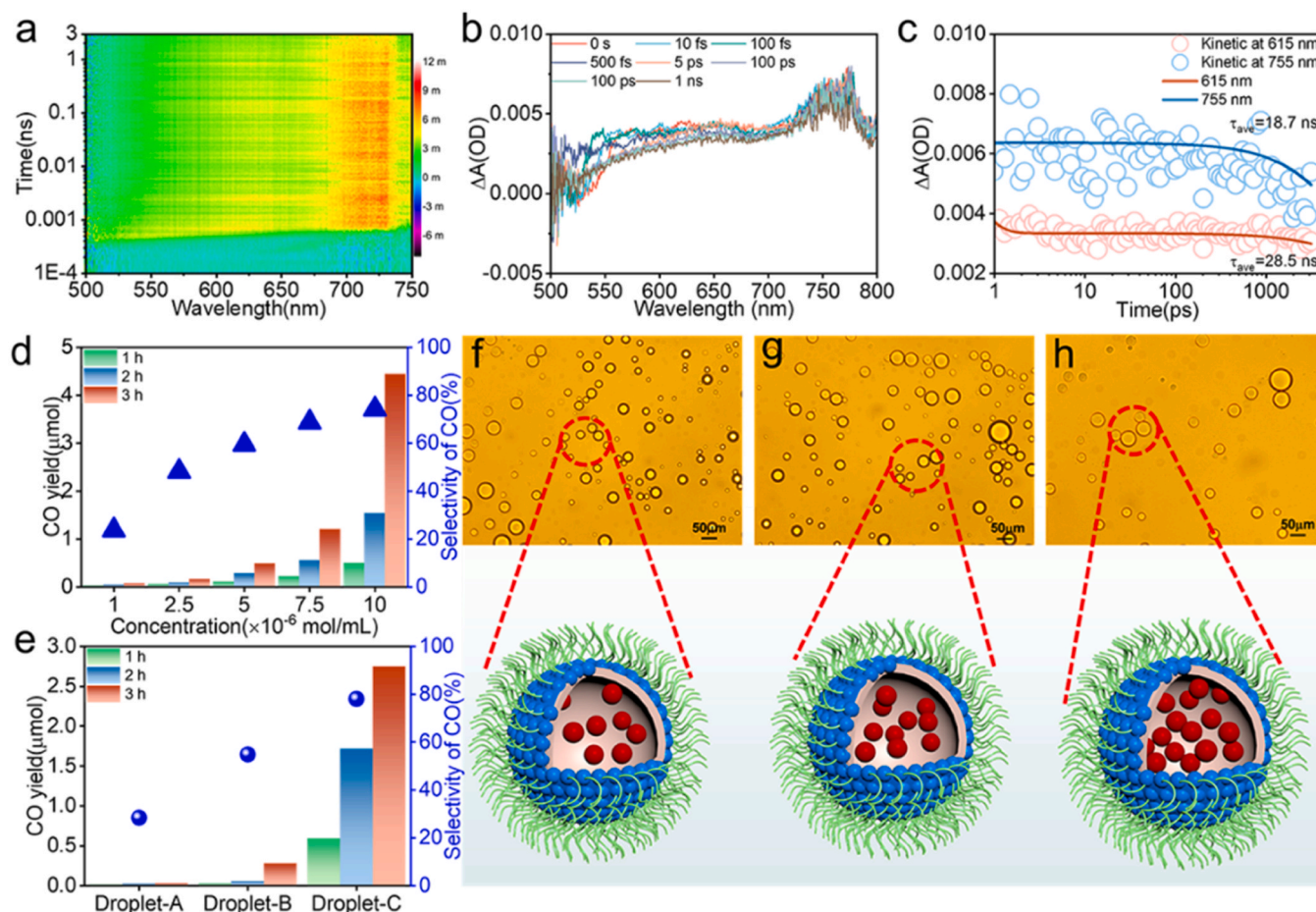

Fig. 4. (a) 2D transient absorption surface plot of Ru(bpy)$_3$Cl$_2$ aggregates, (b) *fs*-TA spectra of the Ru(bpy)$_3$Cl$_2$ aggregates, (c) Decay kinetics probed at 615 and 755 nm, (d) Photocatalytic CO$_2$ reduction activity of Ru(bpy)$_3$Cl$_2$ with different Ru(bpy)$_3$Cl$_2$ concentrations, (e) Photocatalytic CO$_2$ reduction performances of droplet-confined Ru(bpy)$_3$Cl$_2$ under 610 nm light irradiation and (f-h) Optical microscope images of droplet-A, droplet-B, and droplet-C, and schematic illustration of the droplet-confined Ru(bpy)$_3$Cl$_2$ aggregates.

photocatalytic CO$_2$ reduction performance under red light (610 nm) irradiation. As shown in Fig. 4**d**, when the concentration of Ru(bpy)$_3$Cl$_2$ is relatively low (monomer), only trace amounts of CO are detected (3 h

reaction, 0.0734 μmol CO, CO selectivity of 23.7%, H$_2$ as a byproduct). Impressively, as the concentrations of Ru(bpy)$_3$Cl$_2$ increase, the CO yield significantly increases (3 h reaction, 13.2 μmol CO, CO selectivity

of 80.6%). The generated CO was verified from CO_2 reduction through a ^{13}C isotope experiment (Figure S18). Furthermore, the aggregates also display good durability for CO_2 reduction (Figure S19). In addition, some blank experiments were carried out under the same reaction conditions. In the absence of $Ru(bpy)_3Cl_2$ aggregates, without light irradiation, without CO_2 and TEOA, no product was detected (Figure S20). These results indicated that the $Ru(bpy)_3Cl_2$ aggregates and light were essential for CO_2 reduction reaction.

Furthermore, we normalized CO production rates based on the amount of $Ru(bpy)_3Cl_2$. Obviously, with increasing the amounts of $Ru(bpy)_3Cl_2$, CO production rates harshly increase (Figure S21), which indicates that the aggregation of $Ru(bpy)_3Cl_2$ has a major influence on the CO_2 reduction efficiency, and also reveals the existence of strong synergistic effect within $Ru(bpy)_3Cl_2$ aggregates. To further explore the impact of $Ru(bpy)_3Cl_2$ aggregation on the photocatalytic performance toward CO_2 reduction, three water-in-oil droplets containing different concentrations of $Ru(bpy)_3Cl_2$ were prepared. The concentrations of $Ru(bpy)_3Cl_2$ in droplets A, B, and C were 2, 4, and 8 mg/mL, respectively. The total mass of $Ru(bpy)_3Cl_2$ confined in the three droplets is controlled to be 8 mg. Optical microscope images of droplets A, B, and C are presented in Fig. 4f-g. It was found that the rates of CO production are positively correlated with the aggregated degrees of $Ru(bpy)_3Cl_2$ (Fig. 4e). Since the total mass of $Ru(bpy)_3Cl_2$ in the three droplets is the same, the high photocatalytic activity of $Ru(bpy)_3Cl_2$ aggregates confined in droplet C can be safely ascribed to the high aggregated degree of $Ru(bpy)_3Cl_2$. Interestingly, we have also observed that the selectivity of CO increases from 23.7% to 74.2% with increasing the aggregated degrees of $Ru(bpy)_3Cl_2$ (Fig. 4d), consistent with droplet-confined $Ru(bpy)_3Cl_2$ (Fig. 4e). Therefore, the above results undoubtedly confirm that the aggregation of $Ru(bpy)_3Cl_2$ plays a key role in determining the photocatalytic performance of CO_2 reduction.

3.3. Electronic property analysis

MD simulation and DFT calculations confirm that the $Ru(bpy)_3Cl_2$ aggregates exhibit asymmetric structures and large dipole moments. In principle, the dipole polarization effect probably influences the electronic properties of $Ni(bpy)Cl_2$ co-catalyst and thus regulates the activity and selectivity toward CO_2 reduction reaction. To check this, the charge density difference of $Ni(bpy)Cl_2/Ru(bpy)_3Cl_2$ was calculated (Fig. 5a-b). Compared with the $Ni(bpy)Cl_2/Ru(bpy)_3Cl_2$ monomer, the $Ni(bpy)Cl_2/Ru(bpy)_3Cl_2$ aggregates display different charge distributions [56–58]. A Mulliken charge population analysis suggests that the charge transfer amount from the $Ru(bpy)_3Cl_2$ aggregates to the Ni catalytic centers is

higher than that from the $Ru(bpy)_3Cl_2$ monomer to the Ni catalytic center [59,60]. Furthermore, the projected density of states (pDOS) suggests that the d-orbital center of Ni within the $Ni(bpy)Cl_2/Ru(bpy)_3Cl_2$ aggregates up-shifts to the Fermi energy (E_F) level (Fig. 5c) [61–63]. The up-shift of the d-orbital center for Ni probably derives from the strong charge redistribution caused by the dipole polarization effect of aggregated systems. According to the d-band center theory, the higher d-orbital center indicates a stronger binding energy between the Ni catalytic centers and CO_2, which reflects the strengthened activation ability of CO_2 over the Ni catalytic centers, thus promoting the activity and selectivity of CO_2-to-CO reduction [64,65]. Furthermore, the linear sweep voltammetry result also indicates that the aggregation of $Ru(bpy)_3Cl_2$ could indeed promote the catalytic ability of the Ni catalytic center, consistent with the above results and analysis (Figure S22).

In-situ diffuse reflectance infrared Fourier transform spectroscopy (DRIFTS) was performed to monitor the reduction processes of CO_2 over the $Ni(bpy)Cl_2/Ru(bpy)_3Cl_2$ aggregates (Figure S23). The signal at 2355 cm^{-1} corresponds to the symmetric stretching vibration modes of CO_2 [66]. Significantly, the intermediate species of COOH* for CO_2-to-CO (peaks at 1521, 1543 and 1558 cm^{-1}) were witnessed [67]. Furthermore, we calculated the reaction Gibbs free energy change (ΔG) of CO_2-to-CO reduction on the $Ni(bpy)Cl_2/Ru(bpy)_3Cl_2$ aggregates [68, 69]. Optimized geometric structures of various states over the $Ni(bpy)Cl_2/Ru(bpy)_3Cl_2$ aggregates along the reaction pathway are presented in Figure S24. Compared with the $Ni(bpy)Cl_2/Ru(bpy)_3Cl_2$ monomer, the $Ni(bpy)Cl_2/Ru(bpy)_3Cl_2$ aggregates have a lower ΔG for the formation of COOH* species, which is the rate-determining step of CO_2 reduction to CO (Fig. 5d) [70–72]. This result reflects that the aggregation of $Ru(bpy)_3Cl_2$ can lower the energy barrier for the formation of COOH* intermediates. It is widely believed that increased electron density of co-catalysts can boost CO_2 activation because the electrons in co-catalysts could occupy the anti-bonding orbital of CO_2 and force the linear CO_2 to bend. Therefore, the lowered formation energy of COOH* intermediates can be primarily attributed to the high electron density around Ni catalytic centers induced by the dipole polarization effect of $Ru(bpy)_3Cl_2$ aggregates, consistent with the results of Mulliken charge calculation and d-orbital center analysis. To our knowledge, this is the first observation that asymmetric aggregation of $Ru(bpy)_3Cl_2$ light-harvesting unit can modulate the electron structures of Ni co-catalyst and thus boost CO_2 reduction, which will open a new avenue to manipulate the selectivity and activity of CO_2 reduction reaction.

To further improve the activity of CO_2 photoreduction, cyclohexane/acetonitrile mixed solvent was employed to modulate the local interaction within $Ru(bpy)_3Cl_2$ aggregates. MD simulation (Fig. 6a) and $Ru(bpy)_3Cl_2$--$Ru(bpy)_3Cl_2$ contact map (Fig. 6b and S25) suggest that the introduction of cyclohexane has a slight influence on the aggregated process of $Ru(bpy)_3Cl_2$. Interestingly, Hirshfeld surface analysis proposes that cyclohexane influences the intermolecular interactions within $Ru(bpy)_3Cl_2$ aggregates (Fig. 6c and S26) [38,39]. Meanwhile, we calculated the interaction force between $Ru(bpy)_3Cl_2$ and acetonitrile/cyclohexane (Figure S27). The result suggests that the relatively weak pulling effect of cyclohexane can help to promote the $Ru(bpy)_3Cl_2$-$Ru(bpy)_3Cl_2$ intermolecular interactions. Furthermore, we randomly analyzed the local structures of 30 regions based on MD simulation. The statistical analysis reveals that the introduction of cyclohexane can indeed increase the intermolecular interaction strength of neighboring $Ru(bpy)_3Cl_2$ molecules (Figures S28-30). Fig. 6d displays the UV–vis absorption spectra of $Ru(bpy)_3Cl_2$ aggregates after the addition of cyclohexane. Obviously, the optical absorption intensity is improved substantially, signifying a decreased bandgap energy [10,30]. To further check this, we calculated the bandgap values for $Ru(bpy)_3Cl_2$ aggregates with different intermolecular distances. It was found that a closer intermolecular distance of $Ru(bpy)_3Cl_2$ will result in a smaller bandgap (Figure S31), agreeing with the UV–vis absorption result. Besides, the PL lifetime of $Ru(bpy)_3Cl_2$ aggregates increases when cyclohexane is introduced (Fig. 6e and Table S1), signifying an increased
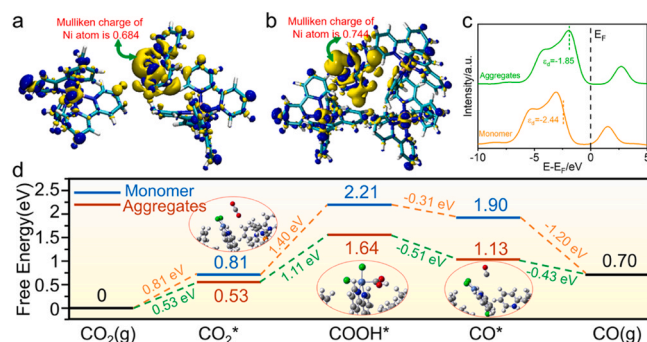

Fig. 5. (a) Charge density difference of $Ni(bpy)Cl_2/Ru(bpy)_3Cl_2$ monomer (isovalue = 0.001 a.u., yellow and blue areas represent charge accumulation and depletion, respectively), (b) Charge density difference of $Ni(bpy)Cl_2/Ru(bpy)_3Cl_2$ aggregates (isovalue = 0.001 a.u., yellow and blue areas represent charge accumulation and depletion, respectively), (c) The pDOS of $Ni(bpy)Cl_2/Ru(bpy)_3Cl_2$ monomer and $Ni(bpy)Cl_2/Ru(bpy)_3Cl_2$ aggregates, and (d) Calculated free energy diagrams for CO_2 reduction over $Ni(bpy)Cl_2/Ru(bpy)_3Cl_2$ monomer and $Ni(bpy)Cl_2/Ru(bpy)_3Cl_2$ aggregates.
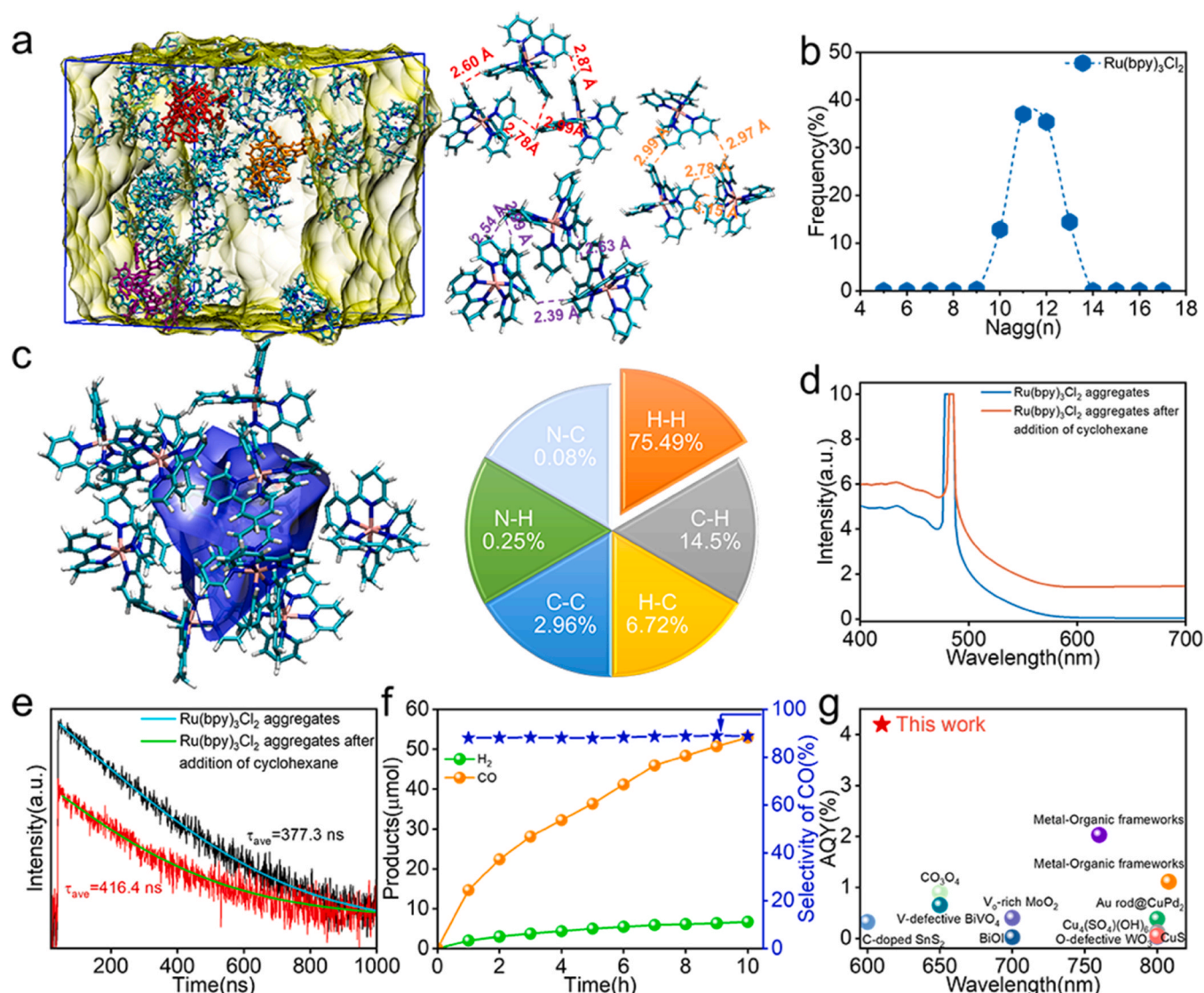
Fig. 6. (a) Structures of Ru(bpy)$_3$Cl$_2$ aggregates in acetonitrile/cyclohexane by MD simulation, (b) The number distribution of Ru(bpy)$_3$Cl$_2$ aggregates based on MD simulation, (c) Hirschfeld surface analysis (mapped over d$_{norm}$), and proportions of intermolecular C···H, H···H, H···C, and other interactions to the total intermolecular interactions based on MD simulated structures, (d) UV–vis absorption spectra of Ru(bpy)$_3$Cl$_2$ aggregates after the addition of cyclohexane, (e) Time-resolved PL spectra of Ru(bpy)$_3$Cl$_2$ aggregates with and without adding cyclohexane, (f) Photocatalytic stability of Ru(bpy)$_3$Cl$_2$ aggregates after adding cyclohexane, and (g) Comparison of CO$_2$ photoreduction performance with state-of-the-art photocatalysts under red/NIR light irradiation (see Table S3 for details).

chance for photogenerated electrons participating in CO$_2$ reduction. As anticipated, the Ru(bpy)$_3$Cl$_2$ aggregates show a 2.6-fold enhancement in CO production with a selectivity of 88% (Fig. 6f and S32). Meanwhile, the Ru(bpy)$_3$Cl$_2$ aggregates exhibit admirable stability (Fig. 6f), well supported by the Cryo-TEM analysis of Ru(bpy)$_3$Cl$_2$ aggregates before and after the reaction (Figure S33) [73,74]. Impressively, the apparent quantum yield (AQY) of Ru(bpy)$_3$Cl$_2$ aggregates for CO production was estimated to be as high as 4.2% at 610 nm, superior to those of the reported red/NIR-light-driven CO$_2$ reduction systems (Fig. 6g and Table S3).

Based on the above analysis, a proposed mechanism of red-light-driven CO$_2$ reduction over the Ru(bpy)$_3$Cl$_2$ aggregates is presented (Fig. 7). Firstly, the Ru(bpy)$_3$Cl$_2$ aggregates capture red light by the unique electronic coupling between adjacent Ru(bpy)$_3$Cl$_2$ molecules and generate electron-hole pairs. Secondly, the dipole polarization induced by the asymmetric aggregation of Ru(bpy)$_3$Cl$_2$ promotes long-range charge delocalization, exciton dissociation, and formation of CT state, accelerating the separation and transfer of photogenerated charge carriers. Importantly, the dipole polarization of Ru(bpy)$_3$Cl$_2$ aggregates
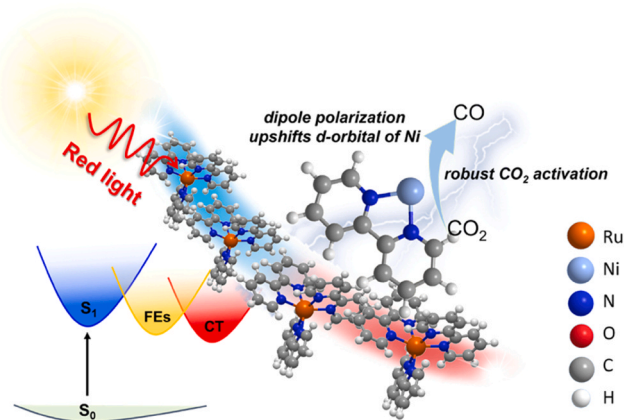

Fig. 7. A proposed mechanism for CO$_2$ reduction over Ru(bpy)$_3$Cl$_2$ aggregates under red light.

is favorable for charge redistribution and hence up-shifts the d-orbital center of Ni co-catalysts, consolidating the activation of CO_2. This photosystem reported here integrates strong red-light harvesting, fast charge separation, and robust CO_2 activation, which synergistically improves the performance toward CO_2 reduction.

4. Conclusion

In this work, we systematically investigated the aggregated effect of $Ru(bpy)_3Cl_2$ on the red-light-driven CO_2 reduction performance. The structures, photophysical properties, and CO_2 photoreduction activity of $Ru(bpy)_3Cl_2$ aggregates were analyzed by various characterizations. Results authorize that the electronic coupling between adjacent Ru $(bpy)_3Cl_2$ molecules induces orbital rearrangements and hybridization, thereby expanding the light absorption edge to the red light region. Meanwhile, the spontaneous symmetry breaking of $Ru(bpy)_3Cl_2$ aggregates accelerates charge separation and transfer, while prolonging the lifetime of charge carriers. Significantly, the symmetry-breaking-induced dipole polarization was found to up-shift the d-orbital center of Ni co-catalysts to the Fermi energy level, which effectively facilitates CO_2 activation and lowers the energy barrier for COOH* intermediates formation. This aggregated photosystem integrates low-energy strong red-light harvesting, fast charge separation, suitable redox potentials, and robust CO_2 activation, which can serve as an unprecedented platform for efficient CO_2 photoreduction. This work deepens the understanding of intermolecular weak forces for tuning the physicochemical properties of aggregates and may offer a way from a new view to control the catalytic activity and selectivity of photocatalytic systems.

CRediT authorship contribution statement

Hongxia Zhang: Funding acquisition, Formal analysis, Data curation. **Hu Shi:** Software, Project administration, Methodology. **Yue Sun:** Methodology, Investigation. **Chun Hao:** Investigation, Data curation. **Pengju Yang:** Writing – review & editing, Writing – original draft, Visualization, Validation, Supervision, Resources, Funding acquisition, Formal analysis, Conceptualization. **Hengquan Yang:** Project administration, Methodology, Investigation. **Jianghong Zhao:** Project administration, Methodology, Investigation.

Declaration of Competing Interest

We declare that we have no financial and personal relationships with other people or organizations that can inappropriately influence our work.

Data availability

No data was used for the research described in the article.

Notes

The authors declare no competing financial interest.

Appendix A. Supporting information

Supplementary data associated with this article can be found in the online version at doi:10.1016/j.apcatb.2024.124222.

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